

# Inverse modeling of basal conditions of the East Antarctic Ice Sheet

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## Background & motivation

- Effect of melting temperate glaciers in the Antarctica is crucial for sea level rise
- Queen Maud Land (QML) and Wilkes Land (WL) considered (see Fig. 1)
- Basal friction plays a major role for flow speed
- Friction coefficient on a glacier's base not observable → Inverse methods:
- Regularization:
  - prevent over-fitting & enhance convexity in inverse methods
  - L-Curve analysis: determine optimal level of regularization

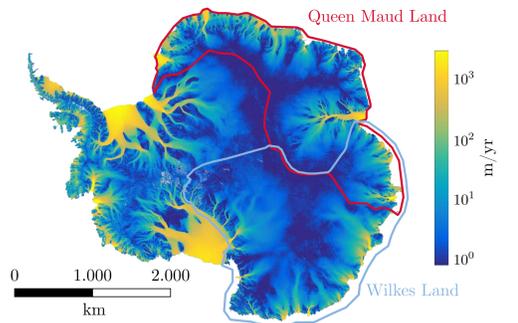


Figure 1: Observed ice flow velocities in the Antarctica [1].

Challenge	Aim
Basal conditions are hard to observe but crucial for the ice flow.	Infer basal drag coefficient and its spatial distribution using an inverse approach.

## Methodology & approach

- NASA's Ice-sheet and Sea-level System Model (ISSM, [2]) used to solve the Higher-Order approximation [3]

$$\frac{\partial}{\partial x_1} \left( 4\eta \frac{\partial u_1}{\partial x_1} + 2\eta \frac{\partial u_2}{\partial x_2} \right) + \frac{\partial}{\partial x_2} \left( \eta \frac{\partial u_1}{\partial x_2} + \eta \frac{\partial u_2}{\partial x_1} \right) + \frac{\partial}{\partial x_3} \left( \eta \frac{\partial u_1}{\partial x_3} \right) = \rho_i g \frac{\partial h_s}{\partial x_1},$$

$$\frac{\partial}{\partial x_1} \left( \eta \frac{\partial u_1}{\partial x_2} + \eta \frac{\partial u_2}{\partial x_1} \right) + \frac{\partial}{\partial x_2} \left( 4\eta \frac{\partial u_2}{\partial x_2} + 2\eta \frac{\partial u_1}{\partial x_1} \right) + \frac{\partial}{\partial x_3} \left( \eta \frac{\partial u_2}{\partial x_3} \right) = \rho_i g \frac{\partial h_s}{\partial x_2},$$

$$\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} = 0.$$

Base boundary condition: linear sliding law

$$(\boldsymbol{\sigma} \cdot \mathbf{n}) \cdot \mathbf{t}_i = \tau_{b,i} = -\alpha^2 u_{b,i} = -k^2 N u_{b,i}, \quad i = 1, 2 \text{ on } \Gamma_{b,gr}$$

- Drag coefficient  $k$  and friction coefficient  $\alpha = N^{1/2} k$  are crucial → quantity of interest
- Effective pressure  $N$ : accounts for the presence of water lubricating the ice-base interface
- Derive the inverse (minimization) problem [4]:

- Interpret  $\mathbf{u}$  as a function of  $k$ ,  $\mathbf{u} = \mathbf{u}(k)$
- Modify base BC  $k$  to fit  $\mathbf{u}(k)$  to  $\mathbf{u}^{obs}$
- Penalize gradient of  $k$  (regularization)

Inverse problem

$$\min_k \frac{1}{2} \int_{\Gamma_s} (u_1(k) - u_1^{obs})^2 + (u_2(k) - u_2^{obs})^2 dx + \lambda \frac{1}{2} \int_{\Gamma_{b,gr}} \|\nabla k\|^2 dx$$

$:= J^{obs}(\mathbf{u}(k))$   
 $:= J^{reg}(k)$

- Tikhonov regularization term  $\lambda > 0$ :
  - More realistic spatial  $k$ -distributions
  - Enhance convexity
  - Determined through L-Curve analysis

## Results & discussion

- Mesh with element sizes of 0.5 km to 64 km used (approx. 350 k elements)
- Characteristic L-shape in the  $(\ln J^{obs}, \ln J^{reg})$ -space
- Regularization cost increases as  $\lambda$  decreases and data cost increases as  $\lambda$  does
- Trade-off curve close to the individual runs (no outliers)

## Queen Maud Land (QML)

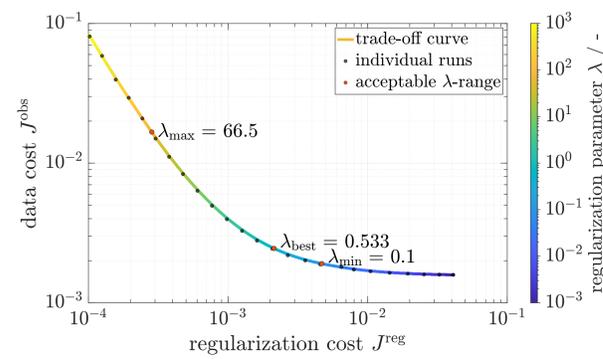


Figure 2: L-Curve results for Queen Maud Land. The red dots identify the best (max. curvature of trade-off), min/max (half of max. curvature of trade-off) value for  $\lambda$  and thus mark the corner region of the curve.

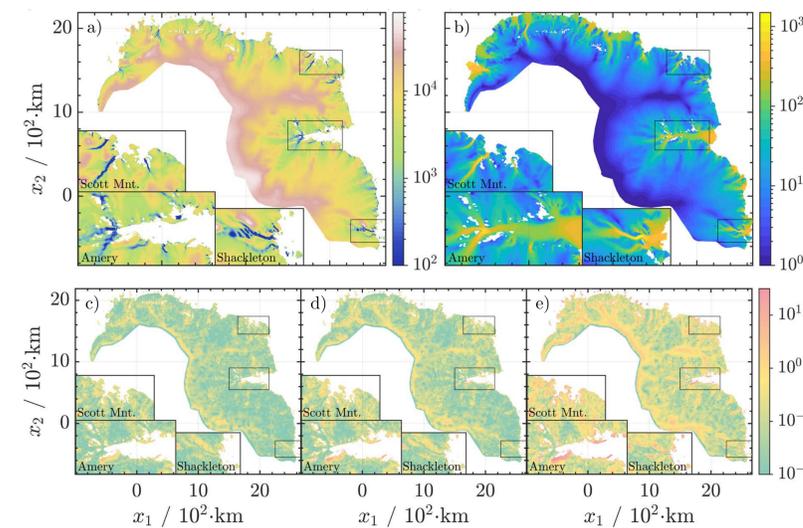


Figure 3: Best estimates for a) the friction coefficient in  $\text{Pa} \cdot (\text{m}/\text{yr})^{-1}$  and b) the velocity in  $\text{m}/\text{yr}$  for QML. Plots c-e) show the relative velocity error for min/best/max values of  $\lambda$  respectively.

- Linear sliding law reflected: fast-flowing regions show small values for  $\alpha$  and inversely (see Fig. 3a)).
- Complex distributions of  $\alpha$  in fast-flowing regions (Amery and Shackleton ice shelves)
- Visual coherence of computed velocity (Fig. 3b) and the observed velocity (Fig. 1)
- Relative velocity misfit (Fig. 3c-e)) unsurprisingly shows smaller errors for smaller  $\lambda$  values and vice-versa
- Overall distribution remains the same across min-/best-/max- $\lambda$

## Wilkes Land (WL)

- Low-res mesh in WL (approx. 120 k elements)
- Higher max- $\lambda$  and larger range as well as overall higher data cost in Fig. 4a) compared to QML
- Results for QML and WL coincide
- Linear sliding law: fast-flowing regions show small values for  $\alpha$  and inversely (see Fig. 4b))

- Convergence issues
  - High-res. Mesh of approx. 350 k elements
  - Lake Vostok was set to grounded area
  - Mountain ranges in Oates Land and Ross Shelf results in badly covered glaciers
  - Cropping of land tongues and islands lead to some improvement

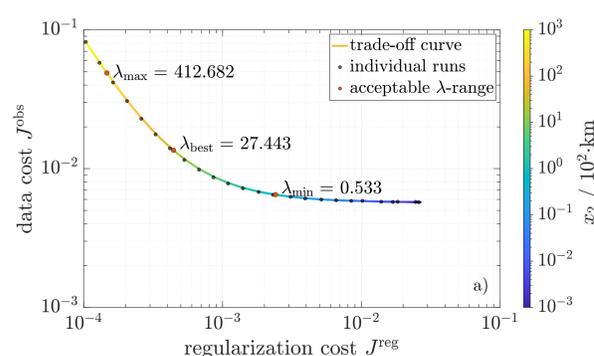
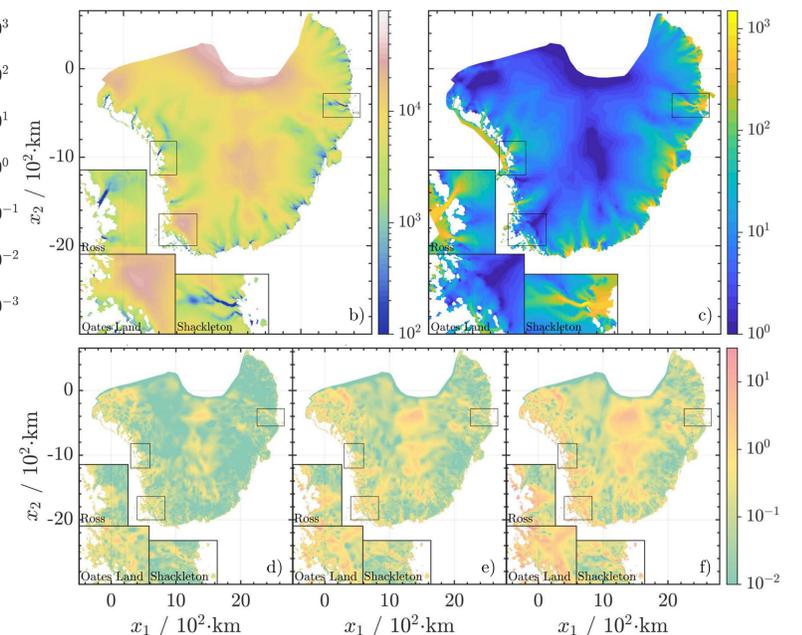


Figure 4: a) L-Curve for Wilkes Land (analogously to Fig. 2). Best estimates for b) the friction coefficient in  $\text{Pa} \cdot (\text{m}/\text{yr})^{-1}$  and c) the velocity in  $\text{m}/\text{yr}$  for Wilkes Land. Plots d-f) show the relative velocity error for min/best/max values of  $\lambda$  respectively.



## Conclusion

- Suitability of Higher-Order approximation and linear sliding law is domain-dependent
- L-Curve approach promising method for identifying the optimal level of regularization
- Optimal distributions for  $k$  and  $\alpha$  could be found
- Overall, results for QML and WL coincide
- Convergence depends on the specific domain

## Future work

- Get convergence for high-res. mesh in Wilkes Land
- Compare results in intersection of QML and WL
- Only the current state in East Antarctica considered (stationary):
  - Derive results for whole Antarctica
  - Expand model for time-dependent problems

Project sea level rise!

## References

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